

Mechanical Vibrations
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Module No #07

Multi DOF

Lecture No. #02

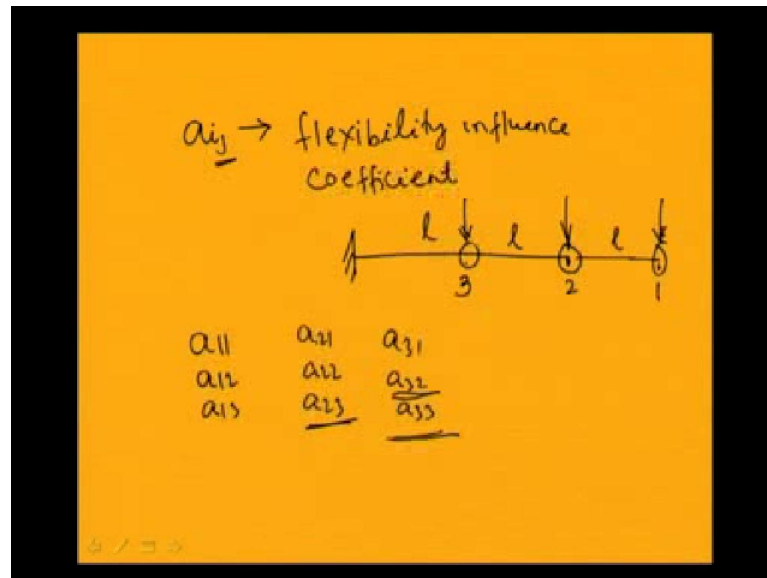
Properties of Vibrating Systems: Flexibility and Stiffness Matrices, Reciprocity Theorem

An IITG person promises only what he can deliver; an IITG person delivers what he promises.

So, today we are going to study about this multi-degree of freedom systems. So, last class we have studied how to determine the equation motion of this multi-degree of freedom system and you know, different methods of finding the equation motion. The methods are either you can use the inertia base principle, that is, Newton's methods or energy based principle, that is, Lagrange principle or Hamilton's principle or extended Hamilton principle.

Also, I have started another, I told you another method to find the equation motion; motion, that is by using the flexible influence coefficient method. And we have seen the definition of flexibility influence coefficient. So, two different types of flexibility influence coefficient, I told, one is the displacement flexibility coefficient and other, it is stiffness flexibility coefficient.

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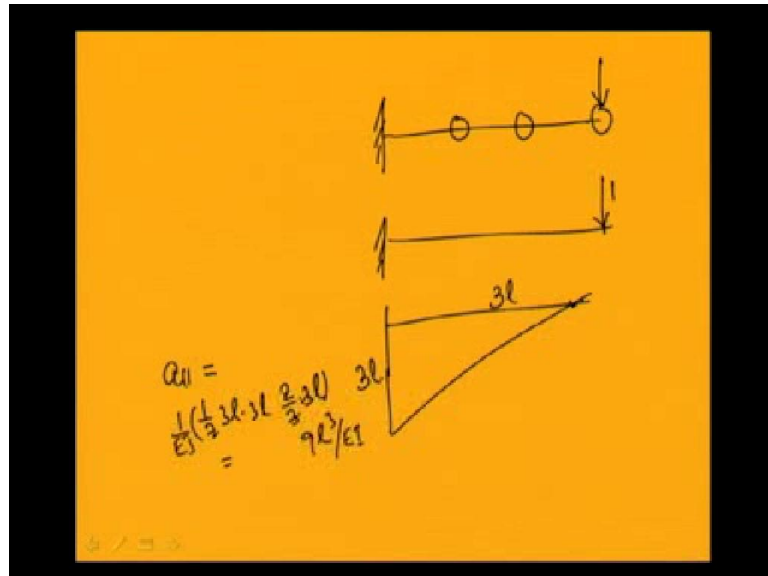


So, that is flexibility influence coefficient, a_{ij} plus, plus, I have defined it as the displacement at i . So, it is the displacement at i due to a unit force applied at j when the other forces are 0. So, a_{ij} is the flexibility influence coefficient. So, this, the displacement flexibility, flexibility, influence coefficient and we have defined these flexibility influence coefficient as the displacement at i due to unit force at j when the forces at other places are 0. Also, we have taken one example to derive or to find the influence coefficient.

So, in that example I have taken three mass, so they are placed at same distance, that is, l , so this is l , l and l . So, I can take the mass, same mass, m_1 , m_2 , m_3 or m . So, I have taken, so this station 1 and this is station 2 and this is station 3 and I have found the influence coefficient a_{11} , a_{12} , a_{13} , a_{21} , a_{22} , a_{23} and a_{31} , a_{32} , a_{33} . So, a_{11} , if the displacement at 1, displacement at 1 due to a unit force at 1; similarly, a_{12} if the displacement at 1 due to a unit force applied at this position, that is, 2 and a_{13} is the displacement at 1 due to a unit force applied at 3. Similarly, a_{21} , a_{21} is displacement at 2. So, this, the displacement at this position, that is, 2 when a unit force is applied at 1. Similarly, a_{22} is the displacement at 2 when a unit force is applied at 2. Similarly, a_{23} is the displacement at 2 when a unit force is applied at 3 and the forces at other places equal to 0. Similarly, a_{31} if the displacement at 3, displacement at 3 when a unit force is applied at 1 and the force at the other two places are 0. Similarly, a_{32} is the displacement at 3 when a unit force is applied at 2 and forces at other places are 0.

Similarly, a 33 is the displacement at 3 when a unit force is applied at 3 and at other forces are 0.

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And I have derived, that those expressions by drawing the Vending Venn diagram. So, displacement of a, so I can consider these as a **contileaver** beam in 3 mass. So, 1, 2 and 3 mass and when a unit force is applied at here, so I can find the displacement. So, I can find the displacement at this position by applying a unit force at this. So, that thing can be obtained by drawing the Vending Venn diagram.

So, the Vending Venn diagram, so here, Vending Venn diagram, you have like this. So, this is the length is $3l$, so this is $3l$. So, to find a 11, one can find a 11 by finding the area of this diagram area. So, area of this diagram equals to half $3l$ into $3l$ and taking the moment of these area. So, moment will be, so it is acting at its length of two-third of $3l$. So, this becomes, so 3, 3 cancel. So, this becomes $9l$ cube by, so this by 1 by EI. So, this becomes $9l$ cube by EI. So, in this way we have determined the influence coefficient a_{11} , so which is equal to $9l$ cube by EI.

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$$a_{21} = \frac{(2l \times l \times l + \frac{1}{2} \times 2l \times 2l \times 2l)}{EI}$$

$$= \frac{2l^3 + 8l^3}{EI} = \frac{10l^3}{EI}$$

(Note: The handwritten result in the image is $\frac{14l^3}{3EI}$)

Similarly, you can determine the other coefficients also. For example, let me determine a 21. So, a 21 is the displacement at 2 due to a unit force at 1 and forces at other place equal to 0. So, when a unit force is applied here, so the displacement can be obtained from this Vending Venn diagram. So, this is the Vending Venn diagram, length is 3l as you are applying a unit force here, vending moment at this position will be 3l. So, the vending moment at, are the displacement at this position, can be obtained from this area, vending moment of this area. So, you can obtain it, vending moment of this area. So, a 21 will be...

So, this area you can divide into 2 parts, so this is the rectangular part and this is the triangular part and a 21 you can obtain from this area as, by taking the moment, if, of this point, for this rectangular part. So, this rectangular part, this length is l. So, this is l and this 2, this length is l. So, this area will be equal to 2l into l into moment. So, you have to take the moment, so this distance is 2l. So, it is at a distance l and for this triangular part it will be half, this is 2l into, so this is 2l and this is also 2l. So, 2l into 2l into, from this to this, centroid of this point is at this, so this is two-third of, two-third of 2l by EI. So, this becomes, so this is 2l cube plus, so 2, 2 cancel, so this becomes, 2 and 2, 4, 28, so 8l cube by 3 by EI. So, this becomes 6 plus 8, this is 14l cube by 3 EI. So, in this way in the last class we have determined the flexibility influence coefficient of all the points.

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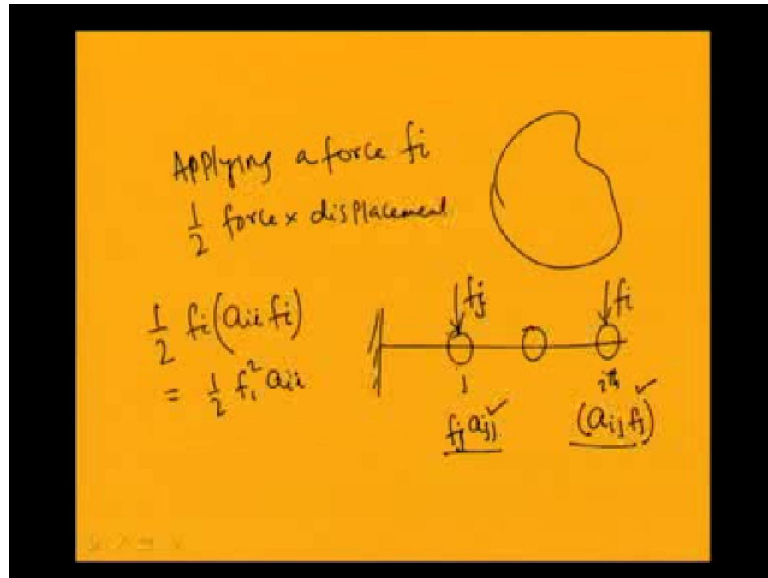
$$a = \frac{l^3}{3EI} \begin{bmatrix} 27 & 14 & 4 \\ 14 & 8 & 2.5 \\ 4 & 2.5 & 1 \end{bmatrix}$$
$$a_{ij} = a_{ji}$$

Reciprocity theorem

$$a_{ij} = a_{ji}$$

So, we have determined the matrix and that matrix is given by, so that matrix already you have seen and that matrix, it is given by l^3 by $3EI$ into 27, 14, 4, 14, 8, 2.5 and 4, 2.5, 1. So, in this case you can observe, that this a_{ij} , a_{ij} equal to a_{ji} . That means, a_{21} , so this a_{21} equal to 14, **your a_{21}** , a_{12} is also 14. Similarly, a_{31} , so this is a_{13} equal to 4 and a_{31} is also 4. Similarly, **a 32** , a_{23} equal to 2.5, a_{32} also equal to 2.5. So, you can see, that are observed, that this a_{ij} equal to a_{ji} . So, this is known as reciprocity theorem, reciprocity theorem. So, for a linear system, so this is a linear system we have considered, so for a linear system you can tell, that this a_{ij} , flexibility influence coefficient a_{ij} equal to a_{ji} . That means, displacement at i due to a unit force at j equal to displacement at j due to a unit force at i . So, to prove that thing let us take a system.

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So, let this is a system or the system what you have, just now you have derived. So, let us take the system and you apply a force, apply a unit or apply a force at i th station. So, let this be i th point applying a force at I , applying a force f_i at i . So, the work done will be equal to half force into displacement.

So, when we are applying a force f_i here, so let this is the f_i force applied here, so that at other place you are not applying any force. So, the displacement at this position will be f_i into displacement will be f_i into a_{ii} as due to a unit force at i th station the displacement will be a_{ii} . So, due to a force f_i applied at i th position, the displacement will be f_i into a_{ii} . And other places at force equal to 0, the work done will be equal to summation of force and displacement at all these places, so as the forces at other places are 0. So, work done will be equal to half into, so this is the force f_i into displacement, f_i into a_{ii} . So, this becomes of, f_i square a_{ii} .

Now, after applying this force f_i lift, again we will apply another force f_j at j th point. So, let this f_j force is acting at j . So, due to this force f_j , this point j will undergo a displacement of f_j into a_{jj} and this i th point, which has already a displacement of a_{ij} , I will have a displacement of a_{ij} , a_{ij} be the displacement at i due to a unit force at j . So, the total displacement at this position due to this force f_j will be a_{ij} into f_j . So, the work done will be due to the displacement at this position and due to the displacement at this position. So, the displacement at this position is a_{ij} f_j and this position, it is f_j into a_{jj} .

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$$\begin{aligned}
 W &= \frac{1}{2} f_i^2 a_{ii} + \frac{1}{2} f_j (f_i a_{ij}) \\
 &\quad + f_i (f_j a_{ij}) \\
 &= \frac{1}{2} f_i^2 a_{ii} + \frac{1}{2} f_j^2 a_{jj} + f_i f_j a_{ij} \quad \text{--- (A)}
 \end{aligned}$$
$$W = \frac{1}{2} f_j^2 a_{jj} + \frac{1}{2} f_i^2 a_{ii} + f_j a_{ji} f_i \quad \text{--- (B)}$$

So, the total work done will be half f_i square a_{ii} plus half, f_j into, f_j into a_{ij} plus half f_i . So, already a force f_i is acting at this position and this point is ongoing a displacement of a_{ij} into f_j . So, due to that, that is the force at, due to that force the work done will be f_i into, $f_i f_j a_{ij}$, f_j into a_{ij} , force f_j into a_{ij} . So, this becomes so the work done w equal to half f_i square a_{ii} plus half f_j square a_{jj} plus, so this point will undergo force into displacement, so this becomes f_i . So, this becomes $f_i f_j a_{ij}$.

Now, let us alternate the application of force, let us first apply the force of f_j , let us first apply the force f_j and then apply the force f_i . So, if we apply the force f_j first, then the work done will be equal to half f_j square a_{jj} and by applying a force f_i after that will have a displacement of, displacement of a_{ji} . So, this is the displacement at j due to a unit force at i and as I am applying a force f_i , so this point will undergo another displacement of a_{ji} into f_i and here a force of f_j is applied to the system. So, this work done will be equal to total work done half f_j square a_{jj} plus half f_i square a_{ii} plus half plus force into displacement, so this in, so this force f_j into a_{ji} into f_i . So, we can, so work done, if we are applying the force at j first and then at i , then the work is half f_j square a_{jj} plus half f_i square a_{ii} and plus half $f_j f_i$ into a_{ji} .

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The image shows a handwritten derivation on a yellow background. The first line is $\frac{1}{2} f_i^2 a_{ii} + \frac{1}{2} f_j^2 a_{jj} + f_i f_j a_{ij}$. The second line is $= \frac{1}{2} f_i^2 a_{ii} + \frac{1}{2} f_j^2 a_{jj} + f_i f_j a_{ji}$. Below this, the equation $a_{ij} = a_{ji}$ is boxed. At the bottom, the text "Stiffness matrix" is underlined.

So, from these two, as the work done in both the cases will be same because it is, it does not depend on the application of the force or the, so we can write, so we can equate these a and b and we can write this half $f_i^2 a_{ii}$ plus half $f_j^2 a_{jj}$ plus half $f_i f_j$ a_{ij} equal to half $f_i^2 a_{ii}$ plus half $f_j^2 a_{jj}$ plus half $f_i f_j$ a_{ji} . So, by equating this you can get this and you can see, that a_{ij} equal to a_{ji} . So, for a , so this is the reciprocity theorem.

And for a linear system you can half a_{ij} , that is, flexibility influence coefficient a_{ij} equal to flexibility influence coefficient a_{ji} , that is, the displacement at i due to a unit force at j equal to a displacement at j due to a unit force at i . So, this is the reciprocity theorem. And now, we will study about the stiffness matrix, stiffness matrix or property of the stiffness matrix of a system.

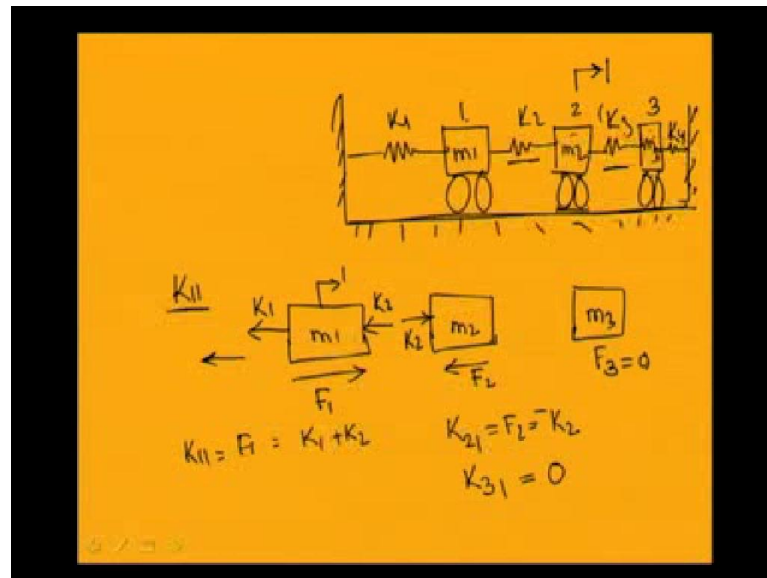
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The image shows a handwritten diagram on a yellow background. At the top, a stiffness matrix K is represented as a 3x3 grid of elements: K_{11} , K_{12} , K_{13} in the first row; K_{21} , K_{22} , K_{23} in the second row; and K_{31} , K_{32} , K_{33} in the third row. Each column is circled, and the columns are labeled K_{1j} , K_{2j} , and K_{3j} respectively. Below the matrix, a definition for K_{11} is given: $K_{11} \rightarrow$ Force required at 1 to have unit displacement while the displacement at other places equal to zero. At the bottom left, the symbol K_{ij} is underlined.

So, in previous system we have written this stiffness matrix K . Let us take a three degree of freedom system. So, the stiffness matrix K can be written as K_{11} , K_{12} , K_{13} , K_{21} , K_{22} , K_{23} and K_{31} , K_{32} , K_{33} . So, this stiffness matrix, this K_{11} or K_{11} can be defined as, so this is the force required at 1, this is the force required at 1 to half, unit displacement at 1. So, this is the force required at 1, so this is the force required at 1 to have unit displacement at 1 while the displacement at other places equal to 0, at other places equal to 0.

So, similarly K_{ij} , you can define as force required at i to have unit displacement at j when the displacement at other places equal to 0, this K_{12} . So, this 1st column, so this 1st column will represent, so this is force required at 1, 2 or 3 to have unit displacement at 1 while displacement at other places equal to 0. Similarly, this K , 2nd column will represent the displacement force required at 1, 2 and 3 to have unit displacement at 2. And similarly, this 3rd column will represent the force required to have unit, this force required at 1, 2 and 3 to have unit displacement at 3 while the displacement at other places equal to 0.

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So, let us take an example. So, let us take a system with three mass and springs and let us determine the stiffness matrix of the system. So, this is K_1 , this is K_2 , this is K_3 and the system is supported and some wheels. **So, the system has...** So, this is mass m_1 , this is mass m_2 and this has mass m_3 .

So, to find K_{11} , that is, force required at 1 to have unit displacement at 1, we can draw the three body diagram of all these three masses and you can find that thing. So, in case of K_{11} , so it is the force required at 1. So, this is 1, this is 2, this is mass 3. So, force required at 1 to have, so this is the force required at 1 to have unit displacement at 1 and displacement at this position and this position, that is at 2 and 3, so could be 0.

So, to draw the three body diagram, let us draw the three body diagram to determine these. So, this is mass 1, mass 2 and this is mass 3. So, we have to find what is the force required at 1 to have unit displacement at 1; so, if it will have the unit displacement at 1. So, this spring will be stretched by $(())$ 1 and so it will exert a force in opposite direction, that is, magnitude will be equal to K_1 . Similarly, the 2nd spring K_2 , so this spring will be compressed by 1 as this point is not moving at, the 2nd point is not moving, so this will be compressed by 1 and so it will exact a force K_2 in this direction.

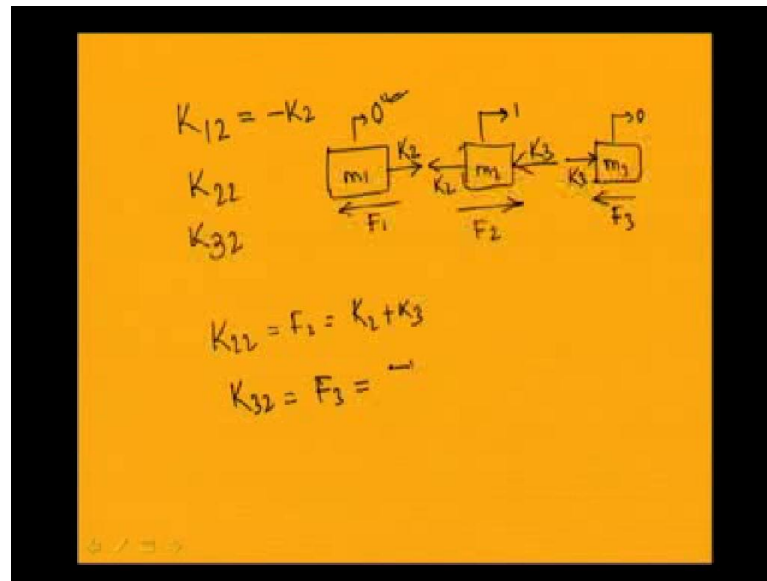
So, for unit displacement of this mass 1, these 2 springs will be giving a force K_1 plus K_2 in this direction. So, they will exact a force in this direction. So, to have that displacement, one has to apply a force F_1 , so that force one has to apply a force F_1

equal to K_1 plus K_2 to have unit displacement at 1 and displacement 0 here and 0 here. So, this K_{11} becomes K_1 plus K_2 . So, K_{11} equal to your F_1 . So, this F_1 equal to K_1 plus K_2 . And now, let us see the 2nd mass, so as the displacement of the 2nd mass equal to 0. So, this spring K_2 , due to this motion, unit motion of mass 1, so it will compress this mass m_2 . So, it will exert a force K_2 and this mass m_2 , and it is K_3 for K_3 spring as there is no motion of mass 2 and 3. So, this K_3 spring will have no motion. So, so it will not exert any force on this. So, the total force acting on m_2 equal to K_2 . So, two half are due to this force.

So, the mass, so moving towards right or to prevent this motion you should apply a force F_2 in this direction, this force F_2 equal to K , so that is equal to K . So, force required at 2 to have unit displacement at 1, displacement at other places equal to 0, so that is K_{21} and so this will be equal to F_2 and it is equal to K_2 . So, this will be equal to minus K_2 because we are taking this, the direction towards right **projectry**, so as we have to apply a force in the opposite direction. So, one has to write, this is equal to minus K_2 .

Similarly, for mass 3, so this is K_3 and this is spring K_4 , so as there is no motion of this mass 2 and 3. So, this spring will, so those spring force acting on the 3 equal to 0 and the spring force acting on 4 equal to 0. So, no force is required as there is no force acting on this mass m_3 , so no force is required to, to have its 0 displacement. So, F_3 will be equal to 0. So, this K_{31} , that is, the force required at 3 to have unit displacement at 1 and displacement at other places equal to 0 will be equal to 0.

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Similarly, you can find K , applying K . So, we can find K_{12} , K_{21} , K_{12} , K_{22} and K_{23} , K_{32} . So, for this, the second mass will have to undergo a displacement of unity and other two mass will be at their equilibrium position. So, they will have no motion. So, for that I can draw the free body diagram.

So, in this case, free body diagram, so mass, so in this mass m_1 will have zero motion and mass m_2 will have unit motion. So, it will move unity and mass m_3 will have also zero motion. So, it will also not move. So, we have to find what is the force required to have this configuration. So, in this case I can draw the free body diagram. So, this, as this mass is not moving or it is stationary, so the spring will not exert the force on this side. But as the second mass is moving, as the spring K_2 , the spring K_2 will be pulled by an amount K , so the spring K_2 will pull this mass m_1 with the force K_2 . So, a force K_2 will act on this mass m_1 . So, to prevent this motion of this mass 1 or to have zero displacement of this mass 1 when force K_2 which acting on this, so we have to apply a force or one has to apply a force F_1 . You have this motion, have 0 motion here, so this F_1 equal to K_{12} , so this will be equal to, so as this is acting in opposite direction, so I can write K_{12} equal to minus K_2 .

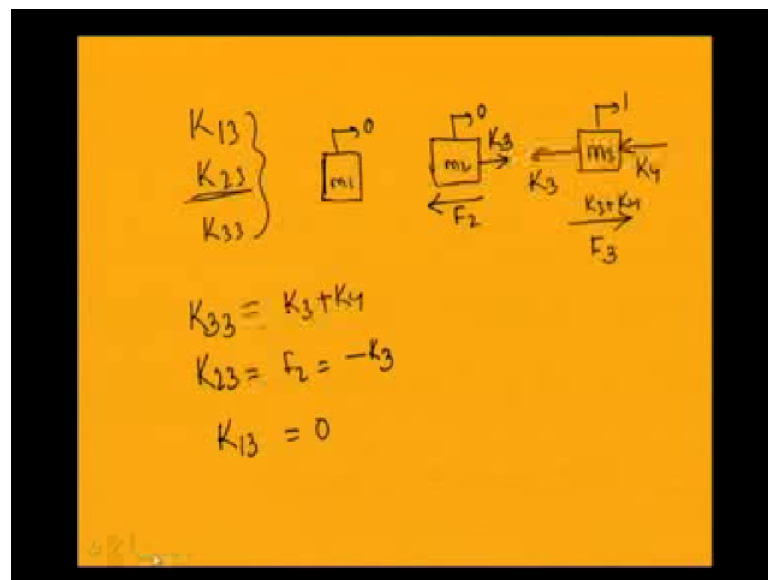
Similarly, when this mass m_2 has unit displacement and mass m_1 and m_3 are zero displacement, this spring K_2 will be pulled by this mass m_2 . So, this spring K_2 will exert a force in opposite direction and so, the free body diagram will be, so it will be K_2

and in the right side. So, the spring K 3, for this spring K 3, the K 3 spring will be compressed by this motion of mass 2 and as there is no motion of mass 3, the relative motion of spring K 3 will be unity. So, the spring K 3 will exert a force in opposite direction and this force will be equal to K 3, K 3 into 1, so that is equal to K 3.

So, a total force of K 2 plus K 3 is acting on mass 2 to, to have unit displacement. A force K 3 and K 2 is acting, so to have that, one has to apply force in this direction. So, that is your F 2 and this force should be equal to K 2 plus K 3 and that is K 22. So, K 22 equal to F 2 equal to K 2 plus K 3. Similarly, this mass 3, so at this time the mass 3 is subjected to a force due to this spring 3, K 3. So, this spring K 3 will push this mass 3 with a force of K 3, as the relative motion of this spring equal to 1 minus 0, that is equal to 1. And as there is no motion of this mass m 3, so K 4 force or force due to this spring K 4 equal to 0. So, the mass 3 will be subjected to a force of K 3 only and that force, so this is K 3.

So, to have this body in equilibrium position, the force required is K 3 in opposite direction, so that is F 3. So, this K 32 equal to F 3 equal to minus K 3.

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Similarly, one can find K 13, K 23 and K 33 by drawing the free body diagram when the mass 3 is having unit displacement and mass 2 and mass 1 having 0 displacement. So, this mass, this is mass m 1 having 0 displacement; mass m 2 having 0 displacement and mass m 3 having unit displacement. So, when mass m 3 has unit displacement, spring K

3 will be pulled by unity. So, it will exert a force in opposite direction, that is, K_3 in this direction. Similarly, the spring 4, K_4 will be compressed by unity and it will exert a force, also in opposite direction, so that is K_4 . So, one has to apply or the force required to have unit displacement at 3 will be equal to K_3 plus K_4 and so that is equal to K , so this is the force F_3 . So, that is equal to K_3 plus K_4 .

So, this force is required to maintain or to have unit displacement at this 1. So, force required at 3, so this is K_{33} , force required at 3 to have unit displacement at 3. So, this is equal to K_3 plus K_4 . Now, to find K_{23} , let us see this mass, free body diagram of this mass m_2 . So, in this case, to have unit displacement at 3, the spring K_3 , so in this case to have unit displacement at 3 plus spring K_3 will pull mass m_2 by m force K_3 in, towards right. So, so this spring K_3 will pull this mass towards right by a force K_3 as there is no motion of mass m_2 and m_1 . So, this force or the spring K_2 will not exert any force. So, the total force acting on mass m_2 equal to K_3 towards right. So, the force equal to have this motion will be in opposite direction to this and this is equal to F_2 and this F_2 equal to K_{23} . So, K_{23} equal to F_2 equal to minus K_3 . Similarly, for mass 1, as there is no motion of this spring K_2 and K_1 , so there is no force acting on mass m_1 , so no force is required to have this motion. So, K_{13} will be equal to 0.

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$$K = \begin{bmatrix} K_1 + K_2 & -K_2 & 0 \\ -K_2 & K_2 + K_3 & -K_3 \\ 0 & -K_3 & K_3 + K_4 \end{bmatrix}$$

$$K_{ij} = K_{ji}$$

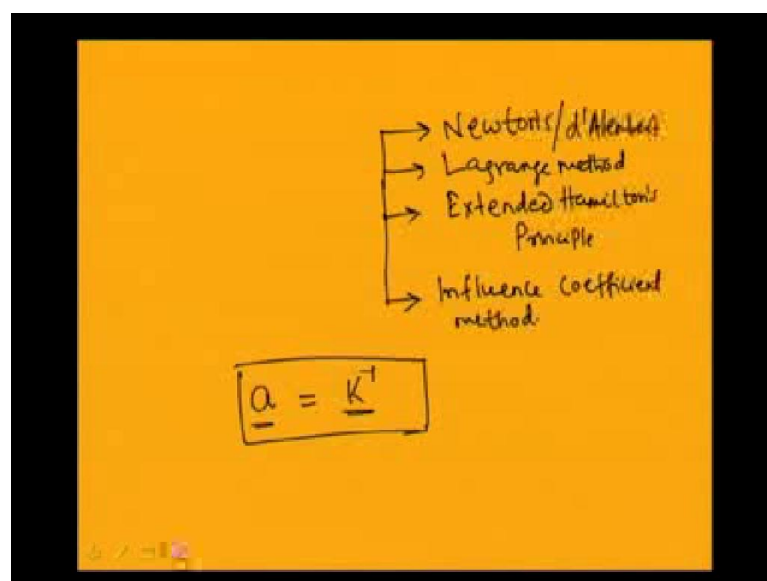
So, in this way we have determined all, the stiffness coefficient. So, by arranging those stiffness coefficients one can write, so this matrix equal to, so this matrix becomes K

matrix equal to $K_1 + K_2 - K_3$ and this is $-K_2 - K_3 + K_4$. So, this is $K_2 + K_3 - K_3$ and this is $0 - K_3 + K_4$. So, here also you can observe, that $K_{ij} = K_{ji}$, so it also follows the reciprocity theorem, that is, $K_{ij} = K_{ji}$, so for a linear system, so you can see, that this stiffness K_{ij} , so $K_{12} = K_{21}$. Similarly, $K_{13} = K_{31}$ and $K_{23} = K_{32}$. So, $K_{ij} = K_{ji}$. So, from the definition of the stiffness influence coefficient, you can determine the influence coefficient or all the elements of this stiffness matrix in this way.

So, in this case, in this example we have, as drawn the free body diagram of these three mass and we have found the force required to have unit displacement to find K_{11} , we have applied a, so to find K_{11} . So, we have what is the force required to have unit displacement at 1 and displacement at other two places equal to 0, and we have found, that is equal to $K_1 + K_2$ and similarly, we have found K_{21} .

So, in this case we have seen that this mass is subjected to a force K_2 to have a unit displacement at 1. So, as this mass is subjected to a force K_{22} , Hamilton equilibrium, the force required at station 2 equal to $-K_2$. Similarly, we have seen, that at this position there is no force acting on this mass 3. So, $K_{31} = 0$. So, in this way, one can determine the elements of this stiffness matrix by finding the force required at station, at i th station, to have force required at i th station to have unit displacement at j th station when the displacement at other places equal to 0.

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Now, till now we have studied how to determine the equation motion of the multi degree of freedom system. So, first, we know, by using the Newton's method we have determined the equation motion Newton's method or d'Alembert principle, and then we have used the Lagrange method and next we have applied extended Hamilton principle to determine the equation motion.

Also, we have studied this influence coefficient method. Now, we have studied this influence coefficient method, coefficient method to determine the mass, to determine the stiffness matrix and the flexibility influence coefficient matrix of the system. So, you can observe, that this displacement flexibility coefficient matrix equal to 1 by K, that is, 1 by this stiffness matrix or reciprocal of the stiffness matrix. So, this is equal to K inverse. So, influence coefficient matrix or displacement influence coefficient matrix or the stiffness influence coefficient matrix, they are related by this formula.

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The image shows a handwritten derivation on a yellow background. The equations are as follows:

$$\begin{aligned}
 & \underline{M\ddot{X} + KX = 0} \\
 & K^{-1}M\ddot{X} + K^{-1}KX = 0 \\
 & \underline{AM\ddot{X} + IX = 0} \\
 & \underline{A\ddot{X} + IX = 0} \quad \leftarrow \\
 & (-A\omega^2 + I)X = 0 \quad \underline{X = X \sin \omega t} \\
 & (A - \frac{1}{\omega^2}I)X = 0
 \end{aligned}$$

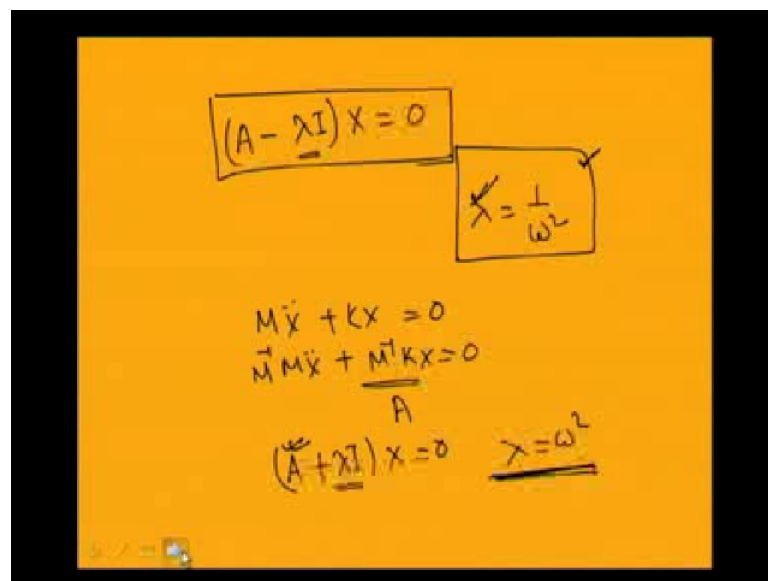
And by finding this displacement coefficient matrix or stiffness matrix one can write the equation motion, by writing this $M \ddot{X} + KX = 0$. So, already we have seen, this is the free vibration or this is the equation motion for free vibration of a multi-degree of freedom system, where m is the mass matrix, K is the stiffness matrix and X is the displacement vector. So, either one can write the equation in this form or one may write this equation by premultiplying K inverse. I can write this equation in this form, so it will be $K^{-1}M \ddot{X} + K^{-1}KX = 0$.

So, $K^{-1}K = I$ and this $K^{-1}K$ inverse equal to A matrix. So, $AM\ddot{X} + I\dot{X} = 0$, so this is $I\dot{X} = 0$.

So, if you know the flexibility influence displacement flexibility influence coefficient, so by multiplying this with mass matrix you can find another matrix, that is, $AM\ddot{X} + I\dot{X} = 0$ proceeding in the previous case to find the normal mode of the system. So, in this case also, we can substitute this \dot{X} equal to ωX , so we can substitute this \dot{X} equal to ωX $I\dot{X} = \omega X$. So, sine ωt or so, by substituting this thing in this equation, so \ddot{X} will become minus $\omega^2 X$.

So, this equation becomes, so I can write this equation, $(A - \omega^2 I)X = 0$. So, in this way one can write this equation and you can, so minus ω^2 plus $I\dot{X} = 0$ or I can write this equation in this form $(A - \omega^2 I)X = 0$.

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Or this equation can be written as $(A - \omega^2 I)X = 0$. So, in this case, in this case, this λ equal to ω^2 . So, if you are finding the Eigen value of matrix A , so this will give you ω^2 , but in the previous case when you have taken the equation in this form $M\ddot{X} + KX = 0$ and you have multiplied it with minus M^{-1} , then you got this equation, $M^{-1}KX = 0$. So, that time also we have written this as A . So, it was reduced to this form, $(A - \omega^2 I)X = 0$.

$\lambda I X$ equal to 0. So, in this case this λ equal to ω square, in the previous case we have seen.

So, the Eigen value of this matrix A equal to ω square and here we were proceeding from the flexibility coefficient, displacement flexibility coefficient, you will get a Eigen value, which is reciprocal of the actual Eigen value of the system. So, this, so in this case the highest frequency will correspond to the lowest Eigen value. And the lowest Eigen frequency will or Eigen value will correspond to the highest Eigen highest frequency of the system. So, you can solve the $(())$, either from this by using this stiffness matrix method or by using this displacement influence coefficient method. In case of displacement influence coefficient method if you proceed, so you will get the Eigen value, which is equal to the reciprocal of this ω square and in this case you can find it, λ equal to ω square.

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Orthogonal Properties of
Eigenvector

$$\underline{X_i} \rightarrow \underline{X_j}$$

$$M\ddot{x} + Kx = 0$$

$$-M\omega^2 x + Kx = 0 \quad \underline{x = X \sin \omega t}$$

$$Kx = M\omega^2 x$$

So, let us see the orthogonal properties of this Eigen vectors; orthogonal Properties of Eigen Vector. So, we have already seen, that this Eigen values corresponding to the, correspond to the square of the natural frequency and Eigen vectors correspond to the normal modes oscillation of the multi-degree of freedom system. So, for the i th mode, so for the i th mode I can write this Eigen vector X_i . If this Eigen vector X_i for the i th mode correspond to the normal mode of the i th mode and X_j will correspond to the normal mode of the j th mode, so one can, so that these, these modes, i th mode and j th

mode are at, these modes are orthogonal with respect to the mass matrix and stiffness matrix of the system.

So, in this case we are assuming that this mass matrix and stiffness matrix are symmetric matrix. So, already we know, that the equation motion can be written in this path and $M \ddot{X} + K X = 0$ by substituting. So, $M \ddot{X} + K X = 0$. I can write this small $M \ddot{X} + K X = 0$. I can substitute this $X = X \sin \omega t$ and I can write this equation minus $M \omega^2 X + K X = 0$. I can write, or capital X also the X can write, minus $m \omega^2 X$ will be, or plus $K X = 0$, or I can write this $K X = m \omega^2 X$.

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$$\begin{aligned}
 K X_i &= \lambda_i M X_i \\
 \lambda_i &= \omega_i^2 \\
 X_j^T K X_i &= X_j^T \lambda_i M X_i \\
 X_j^T K X_i &= \lambda_i (X_j^T M X_i) \quad \text{--- (a)} \\
 K X_j &= \lambda_j M X_j \\
 X_i^T K X_j &= X_i^T \lambda_j M X_j
 \end{aligned}$$

So, for i th mode I can write, for i th mode I can write this $K X_i$ will be equal to $\lambda_i M X_i$, where this λ_i equal to ω_i^2 . So, the previous case, this is ω_i^2 , for the i th mode I can write this equation equal to $K X_i$. So, it will be $K X_i \sin \omega_i t = M \omega_i^2 X_i \sin \omega_i t$. So, this $\sin \omega_i t$ will cancel from both side. So, you can write this $K X_i = M \omega_i^2 X_i$. So, this thing can be written $K X_i = \lambda_i M X_i$. So, here, $\lambda_i = \omega_i^2$.

So, now, premultiplying this transpose of the j th mode. So, I can premultiply this equation by X_j^T . So, this, the transpose of j th mode, so $X_j^T K X_i$ will be equal to $X_j^T \lambda_i M X_i$, or this thing can be written $X_j^T K X_i = \lambda_i X_j^T M X_i$.

into $X_j^T M X_i$. So, let this is equation a. So, let us start from the j th mode. So, I can write for the j th mode, $K X_j$ will be equal to $\lambda_j M X_j$. So, let me premultiply this by X_i^T . So, this $X_i^T K X_j$ will be equal to $X_i^T M X_j$.

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The image shows handwritten mathematical equations on a yellow background. The equations are:

$$X_j^T K X_i = \lambda_j (X_i^T M X_j) \quad \text{--- (b)}$$

$$X_j^T K X_i = \lambda_i (X_j^T M X_i) \quad \text{--- (a)}$$

$$\left. \begin{aligned} X_i^T K X_j &= X_j^T K X_i \\ X_i^T M X_j &= X_j^T M X_i \end{aligned} \right\} \text{For symmetric } M \text{ \& } K$$

$$0 = (\lambda_i - \lambda_j) (X_i^T M X_j)$$

So, I can write this $X_i^T M X_j$. So, this is $X_i^T M X_j$. So, this $X_i^T M X_j$ equal to λ_j into λ_j into $X_i^T M X_j$; $X_i^T M X_j$.

So, now, from this equation a and b, from equation a and b one can write equation, this equation b and equation a is written in this part, $X_j^T K X_i$ equal to $\lambda_i X_j^T M X_i$, we can see this things. So, this is $\lambda_i X_j^T M X_i$, so this is $X_j^T M X_i$. So, this is equation A. Now, the stiffness matrix and mass matrix are symmetric, then you know, that $X_i^T K X_j$ will be equal to $X_j^T K X_i$ and $X_i^T M X_j$ equal to $X_j^T M X_i$ for symmetric. So, this is for symmetric, so this is for symmetric mass and stiffness matrix.

So, from these two by subtracting a from b, so I can write, so if I subtract these things, so becomes, left side becomes 0. So, 0 equal to λ_i minus λ_j , λ_i minus λ_j into $X_i^T M X_j$. So, but this λ_i minus λ_j , that is, this Eigen values are the distinct numbers.

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$$\begin{cases} X_i^T M X_j = 0 \\ X_i^T K X_j = 0 \end{cases} \quad \text{if } i \neq j$$

Orthogonal normal mode
Property

So, λ_i not equal to λ_j and as λ_i not equal to λ_j , then this X_i show, it shows, that $X_i^T M X_j = 0$, if i not equal to j .

So, when i not equal to j , when $X_i^T M X_j = 0$. So, this is the source, the orthogonal property of this Eigen vectors. Similarly, you can show, that $X_i^T K X_j$ also equal to 0. So, the source the orthogonal properties of this mass, orthogonal property Eigen vectors or the normal modes orthogonal properties of this Eigen vector, X_i and X_j or normal mode X_i and X_j , so this orthogonal property of the normal mode, orthogonal property of this normal mode.

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When $i=j$

$$X_i^T M X_j = X_i^T M X_i \rightarrow \text{generalized mass matrix}$$
$$X_i^T K X_j = X_i^T K X_j \rightarrow \text{generalized stiffness matrix}$$

So, when i equal to j . So, this will reduce to, so this will reduce to, so when i equal to j , so when i equal to j , you can see that, so this part equal to 0. So, this will satisfy this equation. So, when i equal to j , this $X_i^T M X_j$ will be equal to $X_i^T M X_i$ and $X_i^T K X_j$ equal to, $K X_j$ equal to $X_i^T K X_j$. So, these are known as, so this is known as generalized mass. So, this will give a diagonal matrix and this diagonal matrix is known as the generalized mass and this is known as the generalized stiffness matrix. This will, if the generalized mass matrix and this will give rise to generalized stiffness matrix.

So, in this way you can obtain the generalized mass matrix and stiffness matrix by using the orthogonal property of this Eigen vector or normal modes.

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$$\underbrace{\begin{bmatrix} 2m & 0 \\ 0 & m \end{bmatrix}}_M \begin{bmatrix} \ddot{x}_1 \\ \ddot{x}_2 \end{bmatrix} + \underbrace{\begin{bmatrix} 3K & -K \\ -K & K \end{bmatrix}}_K \begin{bmatrix} x_1 \\ x_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$
$$\frac{A = M^{-1}K}{\lambda = \omega^2} \qquad \frac{A = K^{-1}M}{\frac{1}{\omega^2}}$$

So, let us take 1 example. So, let us take a system. So, in the system, let 2, m, 0, mass matrix is 2m, 0, 0, m. So, the equation motion is retained in this for 2, m, 0, 0, m, x 1 double dot, x 2 double dot plus 3K, minus K, minus K, K, X 1, X 2, **X 1 X 2**, equal to 0, 0. Let us take this example and let us find, let us find this normal mode and check the orthogonality property of the system.

So, in this case, this is the mass matrix M and this is the stiffness matrix K. So, I can write this A matrix equal to M inverse K. So, either I can find this mass matrix, A matrix of this M inverse K or I can write form K inverse M. So, when I will write A equal to K inverse M, so Eigen value of this will give me 1 by omega square and when I will take A equal to M inverse K, Eigen value lambda will be equal to omega square. So, let me proceed in this way by taking M inverse K.

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$$A = \begin{bmatrix} \frac{3}{2} \frac{K}{m} & -\frac{K}{2m} \\ -\frac{K}{m} & \frac{K}{m} \end{bmatrix}$$
$$|A - \lambda I| = 0, \quad \left. \begin{array}{l} \lambda_1 = \frac{1}{2} \frac{K}{m} \\ \lambda_2 = 2 \frac{K}{m} \end{array} \right\}$$
$$\omega_1 = \sqrt{\frac{K}{2m}}, \quad \omega_2 = \sqrt{\frac{2K}{m}}$$

So, if A equals to M inverse K, so this matrix, again find, so this matrix will become, so A equal to M inverse K and this is equal to 3 by 2 K by M minus K by 2m, this is minus K by M and this is K by M.

So, to, I can find the Eigen value of this matrix. So, Eigen value, so I can find this Eigen value by finding the determinant of A minus lambda I equal to 0. So, this will give me lambda 1 equal to half K by m and lambda 2, lambda, so lambda 2 equal to 2K by m. So, lambda 1 equal to omega 1 square and lambda 2 equal to omega 2 square. So, the natural frequency of the system or the normal mode frequency of the system omega 1 will be equal to root over K by 2m and omega 2 will be equal to root over 2K by m. So, to find the normal mode of this thing, the system, I can find this normal mode by finding the adjoint of A minus lambda I matrix.

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$$X = \text{adj}(A - \lambda I)$$
$$= \begin{bmatrix} \left(\frac{K}{m} - \lambda\right) & \frac{K}{2m} \\ \frac{K}{m} & \frac{3}{2} \frac{K}{m} - \lambda \end{bmatrix}$$
$$\underline{\lambda_i = \lambda_1}$$
$$X_1 = \begin{bmatrix} 0.5 & 0.5 \\ 1.0 & 1.0 \end{bmatrix} \frac{K}{m}$$

So, already you know this, X equal to this normal mode, X is equal to adjoint of A minus lambda i. So, I can find this adjoint matrix. So, adjoint of A minus lambda i becomes K by m minus lambda 1, this is K by 2m and this is K by m and 3 by 2 K by m minus lambda i. So, for lambda i equal to lambda 1, so I can find this X i or X 1, I can find the first normal mode, so X 1, so by substituting that thing and finding the adjoint of this thing.

So, one you can find, this X 1 becomes, so you will get two columns of these and you can check, that by substituting this lambda i equal to lambda 1, so this adjoint matrix, you will get two columns and you can see, that all the, you can verify, that normalize form of this both the columns are same. And this normalized form of both these columns will give this A minus lambda i. You can substitute this lambda i equal to lambda 1. So, you will get this. So, this is equal to 0.1, this is also 0.1, 1.0, 1.0, K by M. So, you can take, that both the columns are same.

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$$x_1 = \begin{bmatrix} 0.5 \\ 1 \end{bmatrix}$$

$$\lambda_1 = \lambda_2 \quad x_2 = \begin{bmatrix} -1 & 0.5 \\ 1 & 0.5 \end{bmatrix} \frac{K}{M}$$

$$x_2 = \begin{bmatrix} -1 \\ 1 \end{bmatrix}$$

So, your x_1 equal to, so x_1 equal to, by normalizing you can write x_1 equal to 0.5, 1. Similarly, by substituting λ_1 equal to λ_2 you can find this x_2 and you can find the adjoint of A minus λ_1 and that will give you minus 1.5, 1.5, K by M or the normalized value of x_2 equal to minus 1, 1.

So, you got the normal modes. So, the first normal mode is x_1 , this is 0.5, 1 and the second normal mode is minus 1, 1. So, after getting these two normal modes you can now verify these normal modes are orthogonal.

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$$x_1^T M x_2 = (0.5 \ 1) \begin{bmatrix} 2m & 0 \\ 0 & m \end{bmatrix} \begin{Bmatrix} -1 \\ 1 \end{Bmatrix}$$

$$= -m + m = 0$$

$$\left. \begin{aligned} x_1^T M x_2 &= 0 \\ x_1^T K x_2 &= 0 \\ x_2^T K x_1 &= 0 \end{aligned} \right\}$$

Now, you just find this X_1 , X_1 dash $M X_2$. So, this will be equal to $0.5, 1$ into $2m, 0, 0,$
 m and $\text{minus } 1, 1$. So, this becomes $\text{minus } M$ plus M equal to 0 . Similarly, you can
 verify, similarly you can verify, that X_2 dash $M X_2$ equal to 0 , X_2 dash $M X_2$ equal to
 X_2 dash $M X_2$ equal to 0 and X_1 dash $K X_2$ equal to 0 and X_2 dash $K X_1$ is also
 equal to 0 . So, this was the orthogonal property of these Eigen vectors or normal modes.

So, today class we have studied about the, about finding the stiffness matrix of a multi
 degree of freedom system and the from, the definition of the stiffness matrix, each
 element of the stiffness matrix, or though K_{ij} element of the stiffness matrix can be
 defined as the force required at i to have unit displacement at j , and displacement at other
 places equal to 0 . And we have seen the reciprocity theorem of the influence coefficient
 and also, we have seen the orthogonal properties of these Eigen vectors or normal
 modes.

So, next class we study the free vibration of, free vibration of multi-degree of freedom
 system by using this modal analysis method.