

Computational Fluid Dynamics for Incompressible Flows
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Lecture 2

Finite difference formulations of Parabolic Equations: Implicit Methods

Hello, everyone so in last class we considered a parabolic equation and we have learned finite difference formulations to discretize this parabolic equations. And mostly we considered in last lecture this explicit method where only one unknown was there. So, in today's lecture, we will consider the same parabolic equation and we will discretize this equation using implicit methods. So, today is lecture two finite difference formulations of parabolic equations and we will learn implicit methods.

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Parabolic Equations

Implicit Method

$$\frac{\partial \phi}{\partial t} = \Gamma \frac{\partial^2 \phi}{\partial x^2}$$

Backward time central space (BTCS) method
Larsson's Method

$$\frac{\phi_i^{n+1} - \phi_i^n}{\Delta t} = \Gamma \frac{\phi_{i+1}^{n+1} - 2\phi_i^{n+1} + \phi_{i-1}^{n+1}}{(\Delta x)^2}$$

$\gamma_x = \frac{\Gamma \Delta t}{(\Delta x)^2}$

$$\phi_i^{n+1} = \phi_i^n + \gamma_x (\phi_{i+1}^{n+1} - 2\phi_i^{n+1} + \phi_{i-1}^{n+1})$$

$$\gamma_x \phi_{i-1}^{n+1} - (1 + 2\gamma_x) \phi_i^{n+1} + \gamma_x \phi_{i+1}^{n+1} = -\phi_i^n$$

tri-diagonal matrix
 unconditionally stable
 $O[(\Delta t), (\Delta x)^2]$

Δt - timestep
 Δx - distance between nodes
 ϕ_i^n at node i , time n
 ϕ_{i-1}^{n+1} , ϕ_i^{n+1} , ϕ_{i+1}^{n+1} at nodes $i-1, i, i+1$, time $n+1$

So, this model governing equation whatever we considered that is $\frac{\partial \phi}{\partial t}$ is equal to $\gamma \frac{\partial^2 \phi}{\partial x^2}$. So, now we will use implicit methods so in implicit method you know that there will be more than one unknowns. Implicit method, the first method what we will learn that is known as Backward Time Central Space method.

So, that is known as BTCS, so Backward Time Central Space, so commonly known as BTCS method. So, that you can see that we will use that finite difference formulation for this first derivative as backward finite difference approximation. And the special second derivative what we have in right hand side that will use central difference method.

So, obviously the order of accuracy will be first order in time and second order in space, so you can see so which we will discretize this equation is so $\frac{\partial \phi}{\partial t}$ will take $\phi_{i,n+1}$, so same thing i is the grid index. And the time level we are denoting with the superscript n or $n+1$ where you know that n is your previous time and $n+1$ is your present or current time.

And difference between this time is known as time step and that is your Δt . So, Δt is your time step, so this now we will use first order accurate scheme which is your backward finite difference formulation for this first derivative with respect to time. So, that we have discretized, now gamma and central difference will use in space.

So, it will be $\phi_{i+1,n+1} - 2\phi_{i,n+1} + \phi_{i-1,n+1}$ divided by Δx^2 . So, as it backward in time all these ϕ will take from the time level $n+1$ which is your present time. So, you can see that we have more than 1 unknowns because at $i+1$ and $i-1$ at these discrete points we have unknown ϕ at time level $n+1$. So, if you see the grid points so you can, so this is your n th level grid and if it is $n+1$ level grid.

So, this is your i this is your $i+1$ and this is your $i-1$. So, this is your obviously you have a uniform grid size so that is your Δx and this is your also Δx . So, it is at level $n+1$ obviously then we will have ϕ_i at $n+1$, ϕ_{i+1} at $n+1$ and ϕ_{i-1} at $n+1$. So, this the time level $n+1$ we are considering and n th level only one known values required. So, that is at point i and its value at ϕ_i in n time level.

So, all the unknown you take in left hand side and obviously there are 3 unknowns so you will get the algebraic equation like so if you define the diffusion factor which is known as $\frac{\gamma \Delta t}{\Delta x^2}$. So, this is your known as diffusion factor coefficient, so $\frac{\gamma \Delta t}{\Delta x^2}$. And now, we can write it as, so we can see $\phi_{i,n+1}$ is equal to $\phi_{i,n} + \frac{\gamma \Delta t}{\Delta x^2} (\phi_{i+1,n} - 2\phi_{i,n} + \phi_{i-1,n})$.

So, all these $n+1$ time level you can take in left hand side so you can see ϕ_i already in the left hand side it is there. So, in the left hand side of these algebraic equation you can see that ϕ_i at time level $n+1$ is there, so this is your term and right hand side also you have $\phi_{i,n+1}$. So, now you can rearrange it and you can write it as $\phi_{i,n+1} - \frac{\gamma \Delta t}{\Delta x^2} (\phi_{i+1,n} - 2\phi_{i,n} + \phi_{i-1,n}) = \phi_{i,n}$ and if you now take $\phi_{i,n+1}$ so this we are taking in the left hand side it will be. So, $\phi_{i,n+1} - \frac{\gamma \Delta t}{\Delta x^2} (\phi_{i+1,n} - 2\phi_{i,n} + \phi_{i-1,n}) - \phi_{i,n} = 0$

ϕ_i at $n+1$. And you have another term plus $\gamma \phi_{i-1}$, sorry $i+1$ at $n+1$.

So, you will have in the then right hand side as ϕ_i will go this side so it will be minus ϕ_i . So, this is the final algebraic equation so at each discrete point you can write these algebraic equations. So, obviously if you form the matrix then you will get a tridiagonal matrix. So, for all the points interior points if you write this in a matrix format then you will get tridiagonal matrix.

Because you can see that you have ϕ_i at $n+1$, ϕ_{i+1} at $n+1$ and ϕ_{i-1} at $n+1$. So, you can see that in the left hand side there are 3 unknowns at $n+1$ time level. So, that is your present time level where we need to find the values at $n+1$ time level and 3 discrete points ϕ_i , ϕ_{i-1} and ϕ_{i+1} .

So, here these 3 discrete points we need to find the value of ϕ . So, obviously you can see it will form tridiagonal matrix and right hand side these value is known from the previous time level n . And obviously you can see this is your coefficient of this ϕ so it is $\gamma \phi_{i-1} + 2\phi_i + \gamma \phi_{i+1}$. And if you can form a tridiagonal matrix considering this equal at all grid points then you will get a tridiagonal matrix and it is easy to solve using TDMA Tri Diagonal Matrix Algorithm which is famously known as Thomas algorithm.

So, if you see the stability analysis of this finite difference formulation then solution is unconditionally stable. So, this is the advantage of this implicit method, so it is unconditionally stable, unconditionally stable. So, that will show later and what is order of accuracy? Order of accuracy is your Δt first order in time and Δx^2 second order in space.

So, this is simplest method where 3 unknowns are there and easy to solve this equation and it is sometime it is known as Laasonen method. So, sometime it is known as Laasonen method. And famously known as BTCS Backward Time Central Space method.

So, obviously you can see as it is implicit method there is no time restriction implicit method and it is unconditionally stable. So, obviously you do not have any time step restriction, so Δt you can choose a larger value but obviously in terms of the accuracy there is some limitation because if you choose higher Δt obviously truncation error also will increase.

Because delta t tends to 0 then only the truncation error will tend to 0. So, obviously you have some limitation in choosing the delta t but you can go larger time step if you use implicit method. But in explicit method we have seen that there is some restriction where gamma x should be less than equal to half for this finite difference formulation of this equation.

So, now we will modify this equation and we will get one finite difference formulation which is second order accurate in time and space and popularly this scheme is known as Crank-Nicolson scheme. It is also implicit scheme and it is known as Crank-Nicolson scheme.

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Parabolic Equations

Crank-Nicolson method

$$\frac{\partial \phi}{\partial t} = \gamma \frac{\partial^2 \phi}{\partial x^2}$$

BTCS

$$\frac{\phi_i^{n+1} - \phi_i^n}{\Delta t} = \gamma \frac{\phi_{i+1}^{n+1} - 2\phi_i^{n+1} + \phi_{i-1}^{n+1}}{(\Delta x)^2}$$

C-N method

$$\frac{\phi_i^{n+1} - \phi_i^n}{\Delta t} = \gamma \frac{\frac{\phi_{i+1}^{n+1} + \phi_{i-1}^{n+1}}{2} - 2\phi_i^{n+1} + \frac{\phi_{i+1}^n + \phi_{i-1}^n}{2}}{(\Delta x)^2}$$

$$\frac{\phi_i^{n+1} - \phi_i^n}{\Delta t} = \frac{\gamma}{2} \left[\frac{\phi_{i+1}^{n+1} - 2\phi_i^{n+1} + \phi_{i-1}^{n+1}}{(\Delta x)^2} + \frac{\phi_{i+1}^n - 2\phi_i^n + \phi_{i-1}^n}{(\Delta x)^2} \right]$$

$O[(\Delta t)^2, (\Delta x)^2]$

BTCS

FTCS

unconditionally stable

So, Crank-Nicolson method or scheme, so here whatever we have used in BTCS scheme let us write first BTCS we have used so what is the governing equation? It is del phi by del t gamma del 2 phi by del x square. So, in BTCS you have used phi i n plus 1 minus phi i n divided by delta t is equal to gamma into phi i plus 1 minus 2 phi i plus phi i minus 1 divided by delta x square. So, this is your just BTCS scheme whatever we have discussed just now and all these you are taking from n plus 1, n plus 1, n plus 1.

So, here what you will do, now instead of taking this diffusion term this dependent variables at n plus 1 time level will take the average of phi i n and phi i n plus 1 that means the values of phi at n th time level and the n plus 1 time level, so the average value you have to take. So, what we will take so these value we will take as phi i plus 1 plus phi i plus 1 at time level n plus 1 and n divided by 2.

So, average value you are going to take. Similarly for this point also you will take $\phi_i + \phi_{i+1}$ divided by 2 and this also we will take $\phi_{i-1} + \phi_i$ divided by 2 at $n+1$ and n by 2. So, now we are modifying this BTCS scheme, in the diffusion term the value of ϕ which we took at $n+1$ time level now we are taking an average value of ϕ_{n+1} and ϕ_n at each discrete points. So, that we are taking.

So, if you write this Crank-Nicolson method which is known also CN method, so if you write it so what you are going to get? $\phi_{i,n+1} - \phi_{i,n}$ divided by Δt is equal to so now you are replacing these as $\frac{\phi_{i+1,n+1} + \phi_{i,n+1}}{2}$ divided by Δx^2 minus $\frac{\phi_{i-1,n} + \phi_{i,n}}{2}$ divided by Δx^2 plus $\frac{\phi_{i-1,n+1} + \phi_{i,n+1}}{2}$ divided by Δx^2 .

So, we have taken the average value, so now you rearrange it so if you rearrange it in this manner then you can see that you will get $\phi_{i,n+1} - \phi_{i,n}$ divided by Δt is equal to. So, you can see that everywhere 2 is there so that you can take it outside gamma by 2 and you take all $n+1$ together so this term, this term, this term and n term you write separately, $\phi_{i,n}$ this and this. So, if you write it then you will get $\phi_{i+1,n+1} - 2\phi_{i,n+1} + \phi_{i-1,n+1}$ and now this term, so this is your $\phi_{i-1,n} + \phi_{i,n}$. So, now divided by Δx^2 .

So, all the $n+1$ time level value we have taken together and similarly you will take all n th time level these terms together. So, you will get $\phi_{i+1,n} - 2\phi_{i,n} + \phi_{i-1,n}$ divided by Δx^2 . So, can you see or observe, can you observe something from this equation. You can see that this term is discretization like BTCS implicit method whatever we have used backward time central space method so similar to that so it is similar to BTCS.

And you can see these all we have taken from the previous time level n which is actually explicit method and we have used FTCS, Forward Time Central Space, so it is kind of Forward Time Central Space method. So, it is you can see Crank-Nicolson method when we are using so you can see when we are using the Crank-Nicolson method it is some average of BTCS and FTCS because right hand side if you see and observe you say see it is average 1 by 2 and it is BTCS type discretization and it is FTCS type discretization.

Why it is BTCS? Because all are at $n+1$ time level, you see this 3 discrete points $i+1, i, i-1$ or at the present time level $n+1$. So, it is kind of BTCS method and this one you can

see all this discrete points $i + 1$, i and $i - 1$ the value of ϕ at n th time level and it is kind of explicit method whatever we learnt FTCS so similar to that. So, and 1 by 2 it is here so you can see it is average of BTCS and FTCS scheme.

The advantage of this discretization is that it is a second ordered time accurate. Anyway we are using central difference method in space so obviously it is Δx square but time level also it is Δt square. So, the accuracy is order of Δt square and Δx square, so second ordered accurate both in time and space and it is unconditionally stable. So, if you do the stability analysis of this finite difference formulation you will find that it is unconditionally stable.

So, you do not have any restriction to choose the time step Δt . So, it is unconditionally stable, so it is unconditionally stable. So, you can see that for this discretization method Crank-Nicolson it is second order time accurate and unconditionally stable. So, how you can show that is a second order time accurate.

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Parabolic Equations

C-N method
Two steps method $\frac{\partial \phi}{\partial t} = \Gamma \frac{\partial^2 \phi}{\partial x^2}$

1st step FTCS

$$\frac{\phi_i^{n+1/2} - \phi_i^n}{\frac{\Delta t}{2}} = \Gamma \frac{\phi_{i+1}^n - 2\phi_i^n + \phi_{i-1}^n}{(\Delta x)^2} \dots (a)$$

2nd step BTCS

$$\frac{\phi_i^{n+1} - \phi_i^{n+1/2}}{\frac{\Delta t}{2}} = \Gamma \frac{\phi_{i+1}^{n+1/2} - 2\phi_i^{n+1/2} + \phi_{i-1}^{n+1/2}}{(\Delta x)^2} \dots (b)$$

Adding equations (a) and (b)

$$\frac{\phi_i^{n+1} - \phi_i^n}{\Delta t} = \Gamma \left[\frac{\phi_{i+1}^{n+1/2} - 2\phi_i^{n+1/2} + \phi_{i-1}^{n+1/2}}{(\Delta x)^2} + \frac{\phi_{i+1}^n - 2\phi_i^n + \phi_{i-1}^n}{(\Delta x)^2} \right]$$

$$\frac{\phi_i^{n+1} - \phi_i^n}{\Delta t} = \frac{\Gamma}{2} \left[\frac{\phi_{i+1}^{n+1/2} - 2\phi_i^{n+1/2} + \phi_{i-1}^{n+1/2}}{(\Delta x)^2} + \frac{\phi_{i+1}^n - 2\phi_i^n + \phi_{i-1}^n}{(\Delta x)^2} \right]$$

$$O[(\Delta t)^2, (\Delta x)^2]$$

So, what method you can see now you can behave it as a 2 step methods, 2 steps method. We are taking about only Crank-Nicolson method, so now you can see as 2 step method let me write the equation first $\frac{\partial \phi}{\partial t}$ is equal to $\Gamma \frac{\partial^2 \phi}{\partial x^2}$. So, this is your n into n plus half and n plus half to n plus 1.

So, you are going in 2 steps from 0 to Δt , so this is your Δt by 2 this is the time step Δt by 2 and this is also when you are going from n plus half to n plus 1 this is also Δt by 2. So,

now you see that this Crank-Nicolson method you can think as a 2 step methods so where you will use first from time level n to n plus half as FTCS which is your explicit method and from n plus half to n plus 1 you will consider as BTCS which is your implicit method.

So, you can see here so you are moving from n to n plus half and the time step is Δt by 2 and you use FTCS method here. And when you are moving from n plus half to n plus 1 so again the time step is Δt by 2 and you can use BTCS, which is your implicit method.

So, if you use that way then you can discretize this model parabolic equation as, so first step so you are using FTCS scheme. So, how we can discretize so you are going from n to n plus half so it is n to n plus half divided by Δt by 2, so which is your time step is equal to now right hand side you have γ so this discretization is ϕ_i plus 1 minus 2 ϕ_i plus ϕ_i minus 1 divided by Δx square.

So, this we are using FTCS so all will be at n th time level. So, if you see the grid points so here you can see so at n to n plus half you are moving so now you use this as point i this is your i plus 1, this is your i minus 1 and constant step size Δx you are using. So, the values of ϕ_i at n you are using ϕ_i plus 1 at n you are using and ϕ_i at n minus 1 you are using and from there you are calculating the value at ϕ_i at n plus half.

So, that from this discretization you are finding. Then in the second step, so in the second step now use BTCS. So, now you are moving from n plus half to n plus 1 so obviously when you are using BTCS the special derivative all you will take from n plus 1 time level. So, it will be ϕ_i at n plus 1 minus ϕ_i at n plus half divided by Δt by 2 is equal to γ ϕ_i plus 1 minus 2 ϕ_i plus ϕ_i minus 1 divided by Δx square and all you will take at n plus 1.

So, all you will take from n plus 1 time level so this is your BTCS method you are using. So, now if you, you can see that if you simply add these 2 equations so if you add this 2 equation what you will get so you see ϕ_i at n plus half, ϕ_i at n plus half you will cancel out. So, if you say that this is your equation a and this is equation b then adding equation a and b what you will get you see? You will get ϕ_i at n plus half so this 2 will get cancelled.

So, this ϕ_i at n plus half and this ϕ_i at n plus half we will get cancelled so you will get ϕ_i at n plus 1 minus ϕ_i at n divided by Δt by 2 is equal to γ , so you are adding it simply we add it so ϕ_i plus 1 at n plus 1 minus 2 ϕ_i at n plus 1 plus ϕ_i minus 1 at n plus 1 divided by Δx

square. And $\phi_{i+1} - 2\phi_i + \phi_{i-1}$ so this you at so these you are taking from $n+1$ because this is your writing so Δx^2 .

So, now this half you can take it this side so it will be $\phi_{i+1} - \phi_i$ divided by Δt so Δt you are taking here and half we are taking in the right hand side. So, you can write $\frac{\gamma}{2}$ and $\phi_{i+1} - 2\phi_i + \phi_{i-1}$ divided by Δx^2 so all at $n+1$ time level plus $\phi_{i+1} - 2\phi_i + \phi_{i-1}$ at n , n , n Δx^2 . So, you can see that you can think this Crank-Nicolson method as a 2 step computations, so one computation you are doing from n to $n+1/2$ as FTCS then $n+1/2$ to $n+1$ as BTCS scheme.

So, you can see that obviously it is unconditionally stable we have already discussed and the order of accuracy is Δt^2 and Δx^2 , so order of accuracy is Δt^2 and Δx^2 . So, we will so in today's lecture that this Crank-Nicolson method is second order accurate in time and space but before that let us discuss another method which is known as beta scheme or beta method.

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Parabolic Equations

Beta method

$$\frac{\phi_i^{n+1} - \phi_i^n}{\Delta t} = \gamma \left[\beta \frac{\phi_{i+1}^{n+1} - 2\phi_i^{n+1} + \phi_{i-1}^{n+1}}{(\Delta x)^2} + (1-\beta) \frac{\phi_{i+1}^n - 2\phi_i^n + \phi_{i-1}^n}{(\Delta x)^2} \right]$$

β - factor

$\beta = 0$ $\frac{\phi_i^{n+1} - \phi_i^n}{\Delta t} = \gamma \frac{\phi_{i+1}^n - 2\phi_i^n + \phi_{i-1}^n}{(\Delta x)^2}$ - Explicit FTCS $O[(\Delta t), (\Delta x)^2]$

$\beta = 1$ $\frac{\phi_i^{n+1} - \phi_i^n}{\Delta t} = \gamma \frac{\phi_{i+1}^{n+1} - 2\phi_i^{n+1} + \phi_{i-1}^{n+1}}{(\Delta x)^2}$ - Implicit BTCS $O[(\Delta t)^2, (\Delta x)^2]$

$\beta = \frac{1}{2}$ $\frac{\phi_i^{n+1} - \phi_i^n}{\Delta t} = \frac{\gamma}{2} \left[\frac{\phi_{i+1}^{n+1} - 2\phi_i^{n+1} + \phi_{i-1}^{n+1}}{(\Delta x)^2} + \frac{\phi_{i+1}^n - 2\phi_i^n + \phi_{i-1}^n}{(\Delta x)^2} \right]$
 $O[(\Delta t)^2, (\Delta x)^2]$ Crank-Nicolson method

For $\frac{1}{2} \leq \beta \leq 1$ the method is unconditionally stable
 For $0 \leq \beta < \frac{1}{2}$ the method is conditionally stable

So, this is a general form of finite difference equation of model equation $\frac{\partial \phi}{\partial t} = \gamma \frac{\partial^2 \phi}{\partial x^2}$ so our governing equation is $\frac{\partial \phi}{\partial t} = \gamma \frac{\partial^2 \phi}{\partial x^2}$. So obviously we have different time level one is n which is

known as previous time level this is your $n + 1$ it is your current time level and the time step is Δt , time step is Δt .

So, we will now use a general form to discretize this equation like this. So, we will discretize as $\phi_{i,n+1} - \phi_{i,n}$ divided by Δt , so the i discrete point obviously if you can see that it will be your i , this is your $i + 1$ and this is your $i - 1$. So, the step size is constant this is your Δx .

Similarly you have here $\phi_{i+1,n}$ which time level n $\phi_{i,n}$ and $\phi_{i-1,n}$ and similarly at $n + 1$ time level also you will have the discrete points $i + 1$ and $i - 1$ so you have $\phi_{i+1,n+1}$, $\phi_{i,n+1}$ and $\phi_{i-1,n+1}$. So, with this now if you discretize this equation using beta method we will write in this so diffusion coefficient γ then we will take a factor beta then we will take the implicit kind of discretization of this special derivative $\phi_{i+1,n} - 2\phi_{i,n} + \phi_{i-1,n}$ divided by Δx^2 .

And we will take this discretization at $n + 1$ time level $n + 1$ and the remaining $1 - \beta$ we will take in n th time level so it will be $\phi_{i+1,n} - 2\phi_{i,n} + \phi_{i-1,n}$ Δx^2 and all these at n th time level. So, what we have done in beta method that we have taken a factor beta where beta times the $\Delta^2 \phi$ by Δx^2 discretization we have taken at the implicit manner.

So, all these ϕ we have taken as n th plus $n + 1$ time level and the remaining $1 - \beta$ times the discretization of this special derivative $\Delta^2 \phi$ by Δx^2 and all these dependent variables ϕ at discrete points $i + 1$ and $i - 1$ we have taken at the time level n . So, that is your previous time level and in the beta times we have taken the value of dependent variable at $n + 1$ time level which is your current time level.

So, now you can see that is a general method we have written in a factor which is beta is a factor. So, now you can see if beta is equal to 1 then what you will get? And if put beta is equal to 0 what you will get? And if beta is equal to half then what you will get? So, all these things already we have discussed.

So, let us say that if you have beta is equal to 0. So, if beta is equal to 0 then you can see the first term here so this first term will get 0 and you will get only the second term and you will get in this form $\phi_{i,n+1} - \phi_{i,n}$ divided by Δt is equal to $\gamma \phi_{i+1,n} - 2\phi_{i,n} + \phi_{i-1,n}$

$\phi_i^n + \phi_{i-1}^n$ divided by Δx^2 . So, what type of discretization it is? So, obviously only 1 unknown is there so it is FTCS, Forward Time and Central Space.

So, explicit method and it is known as FTCS, so now if you put β is equal to 1 special cases we are just discussing. So, β is equal to 1 then obviously you will get $\phi_i^n + 1 - \phi_{i-1}^n$ divided by Δt is equal to $\gamma \phi_i^{n+1} - 2\phi_i^n + \phi_{i-1}^n$ divided by Δx^2 and if β is equal to 1 then this term will get 0 so only this term will remain so you will get $n + 1$.

So, what is this discretization? So, this discretization is obviously implicit method because more than 1 unknown is there and commonly this method is known as BTCS, Backward Time Central Space. So, it is implicit and it is known as BTCS. And if you put β is equal to half then you can see.

So, if you β is equal to half then you are going to get $\phi_i^n + 1 - \phi_{i-1}^n$ divided by Δt is equal to, so β is equal to half $1 - \frac{1}{2}$ will be half so you take common γ by 2, half we have taken common so you will get $\phi_i^{n+1} - \frac{1}{2}(\phi_i^n + \phi_{i-1}^n)$ divided by Δx^2 .

So, now the second term you will get, so $1 - \beta = \frac{1}{2}$ we have taken outside these brackets so $\frac{1}{2}(\phi_{i+1}^n - 2\phi_i^n + \phi_{i-1}^n)$ divided by Δx^2 . So, you know that it is Crank- Nicolson method so this is your Crank- Nicolson method. So, obviously you know that FTCS is conditionally stable and the condition is $\gamma \Delta x$ would be less than equal to half where BTCS is unconditionally stable and also the Crank-Nicolson method is unconditionally stable.

But FTCS and BTCS both are first order accurate in time and second order in space but Crank- Nicolson is second order accurate in both time and space. So, it is order of Δt and Δx^2 this is our order of Δt and Δx^2 but Crank- Nicolson method is order of Δt^2 and Δx^2 .

So, you can see that it is a general method where we have used 1 β factor and with change of β you can get different scheme which may be explicit or implicit or first order accurate or second order accurate depending on the value of β . So, in general we can write for half less

than equal to beta less than equal to 1 the method is unconditionally stable, unconditionally stable.

But if beta varies between less than half and greater than equal to 0 then the method is conditionally stable. So, now will show that Crank- Nicolson method is second order accurate in time and space so how will sum? Because from whatever discretization we have written here from there you cannot easily tell that it is second order or accurate in time. Obviously space you can see that space you can say that it is second order accurate in space but from this discretization of this time derivative easily you cannot say that it is a second order accurate in time.

So, let us derive the truncation error and from truncation error we will be able to find that what the order of accuracy, so truncation error obviously you know the difference between the partial differential equation and the finite difference equation. So, now let us see it.

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Parabolic Equations

Crank-Nicolson method

PDE $\frac{\partial \phi}{\partial t} = \Gamma \frac{\partial^2 \phi}{\partial x^2}$

FDE $\frac{\phi_i^{n+1} - \phi_i^n}{\Delta t} = \frac{\Gamma}{2} \left[\frac{\phi_{i+1}^{n+1} - 2\phi_i^{n+1} + \phi_{i-1}^{n+1}}{(\Delta x)^2} + \frac{\phi_{i+1}^n - 2\phi_i^n + \phi_{i-1}^n}{(\Delta x)^2} \right]$

Taylor series expansion

$\phi_{i+1}^{n+1} = \phi_i^n + (\Delta x \frac{\partial}{\partial x} + \Delta t \frac{\partial}{\partial t}) \phi_i^n + \frac{1}{2!} (\Delta x \frac{\partial}{\partial x} + \Delta t \frac{\partial}{\partial t})^2 \phi_i^n + \text{HOT}$

$\phi_{i-1}^{n+1} = \phi_i^n + \Delta t \frac{\partial \phi}{\partial t} \Big|_i + \frac{(\Delta x)^2}{2!} \frac{\partial^2 \phi}{\partial x^2} \Big|_i + \frac{(\Delta t)^3}{3!} \frac{\partial^3 \phi}{\partial t^3} \Big|_i + \text{HOT}$

$\phi_{i+1}^n = \phi_i^n + \Delta x \frac{\partial \phi}{\partial x} \Big|_i + \frac{(\Delta x)^2}{2!} \frac{\partial^2 \phi}{\partial x^2} \Big|_i + \frac{(\Delta x)^3}{3!} \frac{\partial^3 \phi}{\partial x^3} \Big|_i + \text{HOT}$

$\phi_{i-1}^n = \phi_i^n - \Delta x \frac{\partial \phi}{\partial x} \Big|_i + \frac{(\Delta x)^2}{2!} \frac{\partial^2 \phi}{\partial x^2} \Big|_i - \frac{(\Delta x)^3}{3!} \frac{\partial^3 \phi}{\partial x^3} \Big|_i + \text{HOT}$

LHS of FDE

$\text{LHS} = \frac{\phi_i^{n+1} - \phi_i^n}{\Delta t} = \frac{\partial \phi}{\partial t} + \frac{\Delta t}{2} \frac{\partial^2 \phi}{\partial t^2} + \frac{(\Delta t)^2}{6} \frac{\partial^3 \phi}{\partial t^3} + \text{HOT}$

Term I of RHS

$\frac{\Gamma}{2(\Delta x)^2} [\phi_{i+1}^{n+1} - 2\phi_i^{n+1} + \phi_{i-1}^{n+1}] = \frac{\Gamma}{2(\Delta x)^2} \left[\phi_i^n + \Delta x \frac{\partial \phi}{\partial x} + \frac{(\Delta x)^2}{2} \frac{\partial^2 \phi}{\partial x^2} + \frac{(\Delta x)^3}{6} \frac{\partial^3 \phi}{\partial x^3} + \frac{(\Delta x)^4}{24} \frac{\partial^4 \phi}{\partial x^4} + \dots - 2\phi_i^n \right]$

$= \frac{\Gamma}{2} \left[\frac{\partial^2 \phi}{\partial x^2} + \frac{(\Delta x)^2}{12} \frac{\partial^4 \phi}{\partial x^4} + \dots \right]$

So, what is the discretization of Crank- Nicolson method we have done? So, first let us write that, so the equation is del phi by del t so this is a gamma del 2 phi by del x square. So, this is your partial differential equation and what is your finite difference equation so after discretization using Crank- Nicolson method whatever you get that will be finite difference equation so that will be phi i n plus 1 minus phi i n divided by delta t is equal to gamma by 2.

So, $\phi_{i+1} - 2\phi_i + \phi_{i-1}$ divided by Δx^2 and at time levels let us say $n+1$ and you have $\phi_{i+1} - 2\phi_i + \phi_{i-1}$ this is your n, n, n . So, this is your Crank- Nicolson method. Already we have discussed in detail so now we want to find what is the truncation error?

For that what we will do each dependent variable now we will expand using Taylor series. Then we will see the difference between these partial differential equation and the finite difference equation and then we will find the truncation error and that will tell us that what is the accuracy of this scheme.

So, now you see if you say $\phi_{i+1} - 2\phi_i + \phi_{i-1}$. So, if you use Taylor series expansions, of each dependent variable at particular discrete point and particular at time level. So, you can see this is having $i+1$ and $n+1$ and that means you have $x + \Delta x$ and $t + \Delta t$, so it is function of x and t both.

So, $x + \Delta x$ and $t + \Delta t$ so now if you expand so what you will get? So you will get $\phi_{i,n}$ plus you can write $\frac{\Delta x}{2!} \frac{\partial^2 \phi}{\partial x^2} + \frac{\Delta t}{2!} \frac{\partial^2 \phi}{\partial t^2}$ obviously at level i and n then you can write 1 by factorial 2 . So, then it will be $\frac{\Delta x}{2} \frac{\partial^2 \phi}{\partial x^2} + \frac{\Delta t}{2} \frac{\partial^2 \phi}{\partial t^2}$ square then $\phi_{i,n}$ and other higher order terms.

So, we are not writing other terms so similarly you do for other so now you write $\phi_{i,n+1}$ so only it is $t + \Delta t$ and X_0 it will be simple $\phi_{i,n} + \Delta t \frac{\partial \phi}{\partial t}$ at time level discrete point i and n then plus $\frac{\Delta t^2}{2!} \frac{\partial^2 \phi}{\partial t^2}$. Then you will get $\frac{\Delta t^3}{3!} \frac{\partial^3 \phi}{\partial t^3}$ and other high order terms.

So, that we are not going to write and similarly you write $\phi_{i-1,n}$ and $n+1$. Similarly you write $\phi_{i-1,n+1}$, so you will get $\phi_{i,n}$. So, it is for $\phi_{i-1,n}$ so it is $x - \Delta x$ and $n+1$ is $t + \Delta t$, so about x and t you are expanding. So, obviously you will get minus Δx because you have $\Delta x - \Delta x$ so $\frac{\Delta x}{2!} \frac{\partial^2 \phi}{\partial x^2} + \frac{\Delta t}{2!} \frac{\partial^2 \phi}{\partial t^2}$ it is $t + \Delta t$ only so $\phi_{i,n} + 1$ by factorial 2 so it will be minus $\Delta x \frac{\partial^2 \phi}{\partial x^2} + \frac{\Delta t}{2} \frac{\partial^2 \phi}{\partial t^2}$ plus high order term.

So, now the next will expand $\phi_{i+1,n}$. So, whatever are there this dependent variables so $\phi_{i+1,n}$. So, here $x + \Delta x$ and t so $x + \Delta x$ only so it will be $\phi_{i,n} + \Delta x \frac{\partial \phi}{\partial x}$

ϕ by Δx i n plus Δx square by factorial 2 $\Delta^2 \phi$ by Δx square. Then plus Δx cube by factorial 3 $\Delta^3 \phi$ by Δx cube plus high order term.

Similarly, now we will write ϕ i minus 1 and n . So, now it is minus Δx , x minus Δx so it will be ϕ i n minus Δx $\Delta \phi$ by Δx plus Δx square by factorial 2 $\Delta^2 \phi$ by Δx square i n minus Δx cube by factorial 3 $\Delta^3 \phi$ by Δx cube i n plus high order term.

So, now we have expanded all these ϕ i plus 1, n plus 1 and whatever dependent variables are coming so we have expanded using Taylor series expansion, now let us put it in the finite difference equation. So first, let us consider only the left hand side, if you consider the left hand side then what is there in the left hand side of the finite difference equation.

So, what you will get? Left hand side is ϕ i n plus 1 minus ϕ i n divided by Δt . So, what you can write from the second so from the second directly you can write you can see from here. So, from here you can directly write ϕ i n plus 1 minus ϕ i n divided by Δt so it will be $\Delta \phi$ by Δt , we are not going to write i and n because all that derivatives will be at the point discrete point i and the time level n .

So, plus so Δt we have divided so it will be Δt by 2 Δt by 2 $\Delta^2 \phi$ by Δt square. Then plus you can see from there so Δt you have divided so it will be Δt square divided by 6 factorial 3 means 6 so $\Delta^3 \phi$ by Δt cube plus high order term. So in left hand side we have already seen that only these terms will be there and the right hand side you see so let us say this is your let us say this term this term is 1 and this term is 2 .

So, now you see in the right hand side let us say the term 1, term 1 of right hand side. So, what you are going to get? So, it is γ by 2 ϕ i plus 1 and you can write Δx square also, so you can write γ by 2 Δx square and you have ϕ i plus 1 minus 2 ϕ i plus ϕ i minus 1 in whose time level in n .

So, term 1 let us see what will happen so you take ϕ i plus 1 n so this one then 2 ϕ i anywhere it is there and ϕ i minus 1 n , so these two last two this expansion you take and put it here so what you are going to get here? So, you are going to get here γ by 2 Δx square so ϕ i plus 1 n so already you have expanded using Taylor series so you put it there.

So, it will be $\phi_i^n + \Delta x \frac{\partial \phi}{\partial x} + \frac{\Delta x^2}{2} \frac{\partial^2 \phi}{\partial x^2} + \frac{\Delta x^3}{6} \frac{\partial^3 \phi}{\partial x^3} + \dots$ and other terms high order terms we are not writing then minus, so minus now $2 \phi_i$ so minus $2 \phi_i$ so minus $2 \phi_i$ then you write ϕ_i^{n-1} the expansion whatever you have written.

So, it will be plus $\phi_i^n - \Delta x \frac{\partial \phi}{\partial x}$, all these derivative about the point i and n so about discrete point i and the time level n . So, $\frac{\Delta x^2}{2} \frac{\partial^2 \phi}{\partial x^2}$ then it will be minus $\frac{\Delta x^3}{6} \frac{\partial^3 \phi}{\partial x^3}$. So, it will be 6 so it will be $\frac{\Delta x^4}{24} \frac{\partial^4 \phi}{\partial x^4}$ plus high order term.

And here also you write another term, so write $\frac{\Delta x^4}{24} \frac{\partial^4 \phi}{\partial x^4}$ the other term and minus $2 \phi_i^n$. So, now let us see so you can see that this is your ϕ_i^n this is your ϕ_i^n and it is minus $2 \phi_i^n$ so it is cancelled out. These will cancelled out these also will cancelled out plus and minus.

So, whatever terms you have let us write down here so you will get so Δx^2 so Δx^2 you just divide with each term so you are going to get gamma by 2 . So, gamma by 2 so you can see here $\frac{\Delta x^2}{2} \frac{\partial^2 \phi}{\partial x^2}$ is there so 2 and divided by 2 is there so it will become and Δx^2 will be cancelled out so you will get $\frac{\partial^2 \phi}{\partial x^2}$.

So, next you can see this term so you have 2 into Δx^4 by $24 \frac{\partial^4 \phi}{\partial x^4}$ to the power 4 . So, if you 2 into this you write then you will get and Δx^2 will be if you divide then it will be Δx^2 so it will be Δx^2 divided by 2 , 2 was there so it will be 12 , $\frac{\partial^4 \phi}{\partial x^4}$. And plus other high order terms so that we are not writing.

So, now in the first term in the right hand side we have done similarly let us do the second term in the right hand side. So, let us do the second term, so in the second term you can see this is ϕ_i^{n+1} all n plus 1 .

(Refer Slide Time: 51:15)

Parabolic Equations

Term 11 of RHS

$$\frac{\Gamma}{2(\Delta x)^2} [\phi_{i+1}^{n+1} - 2\phi_i^{n+1} + \phi_{i-1}^{n+1}]$$

$$= \frac{\Gamma}{2(\Delta x)^2} \left[\phi_i^n + \Delta x \frac{\partial \phi}{\partial x} + \Delta t \frac{\partial \phi}{\partial t} + \frac{(\Delta x)^2}{2} \frac{\partial^2 \phi}{\partial x^2} + (\Delta x)(\Delta t) \frac{\partial^2 \phi}{\partial x \partial t} + \frac{(\Delta t)^2}{2} \frac{\partial^2 \phi}{\partial t^2} \right.$$

$$+ \frac{(\Delta x)^3}{6} \frac{\partial^3 \phi}{\partial x^3} + \frac{(\Delta t)^3}{6} \frac{\partial^3 \phi}{\partial t^3} + 3 \frac{(\Delta x)^2 \Delta t}{6} \frac{\partial^3 \phi}{\partial x^2 \partial t} + 3 \frac{(\Delta t)^2 \Delta x}{6} \frac{\partial^3 \phi}{\partial x \partial t^2} + \dots$$

$$- 2\phi_i^n - 2\Delta x \frac{\partial \phi}{\partial x} - 2 \frac{(\Delta t)^2}{2} \frac{\partial^2 \phi}{\partial t^2} - \frac{2(\Delta t)^3}{6} \frac{\partial^3 \phi}{\partial t^3} - 2 \frac{(\Delta t)^4}{24} \frac{\partial^4 \phi}{\partial t^4} + \dots$$

$$+ \phi_i^n - \Delta x \frac{\partial \phi}{\partial x} + \Delta t \frac{\partial \phi}{\partial t} + \frac{(\Delta x)^2}{2} \frac{\partial^2 \phi}{\partial x^2} - \Delta x \Delta t \frac{\partial^2 \phi}{\partial x \partial t} + \frac{(\Delta t)^2}{2} \frac{\partial^2 \phi}{\partial t^2}$$

$$- \frac{(\Delta x)^3}{6} \frac{\partial^3 \phi}{\partial x^3} + \frac{(\Delta t)^3}{6} \frac{\partial^3 \phi}{\partial t^3} + \frac{3(\Delta x)^2 \Delta t}{6} \frac{\partial^3 \phi}{\partial x^2 \partial t} - \frac{3(\Delta x)(\Delta t)^2}{6} \frac{\partial^3 \phi}{\partial x \partial t^2} + \dots \Big]$$

$$= \frac{\Gamma}{2(\Delta x)^2} \left[2 \frac{(\Delta x)^2}{2} \frac{\partial^2 \phi}{\partial x^2} + 2 \times 3 \frac{(\Delta x)^2 \Delta t}{6} \frac{\partial^3 \phi}{\partial x^2 \partial t} + \dots \right]$$

$$= \frac{\Gamma}{2} \left[\frac{\partial^2 \phi}{\partial x^2} + \Delta t \frac{\partial^3 \phi}{\partial x^2 \partial t} + \dots \right]$$

So, you can write term 2 of right hand side so it will be gamma by 2 and delta x square if you take outside. Then it will be delta x square and all at n plus 1 level, phi i plus 1 minus 2 phi i plus phi i minus 1, n plus 1, n plus 1, n plus 1. So, now whatever in earlier slide we have shown that this and this and this so these expansion you put it there.

So, if you put it there what you going to get? So, you are going to get equal to so it will be gamma by 2 delta x square I am just directly putting all these expansion in last slide whatever we have shown. So, you will get phi i n plus del x del phi by del x about point i and time level n I am not going to write this plus del t del phi by del t plus del x square by factorial 2 is 2.

Then del 2 phi by del x square then you write del x into del t. So, it will be del 2 phi by del x del t plus del t square by factorial 2 means 2 it will be del 2 phi by del t square then another cube is there. So, that if you write it here you will get del x cube by factorial 3, so factorial 3 is 6 you can write directly.

Del cube phi by del x cube then plus del t cube by factorial 3 means 6 del cube phi by del t cube then you will get 3 into del x square delta t divided by factorial 3 so it will be 6 it will be 6. Then del cube phi by del x square del t plus 3 into del t square del x by factorial 3 means 6 you write and del cube phi by del x del t square and other high order terms. So, now these we have written the expansion for phi i plus 1 n plus 1.

Now, the term is $\frac{1}{2} \phi^{i+n} + 1$ so that we are going to write the expansion so it will be multiplied by 2, because $\frac{1}{2} \phi^{i+n}$ is there. So, whatever you got so $\frac{1}{2} \phi^{i+n} - 2 \Delta t \frac{d\phi}{dt} + \frac{2 \Delta t^2}{2!} \frac{d^2\phi}{dt^2} - \frac{2 \Delta t^3}{3!} \frac{d^3\phi}{dt^3} + \frac{2 \Delta t^4}{4!} \frac{d^4\phi}{dt^4} + \dots$ So, and other high order term.

So, this is the expansion for $\phi^{i+n} + 1$. Now, let us write the expansion of $\phi^{i-n} + 1$ at time level $n + 1$. So, now let us write the expansion of $\phi^{i-n} + 1$. So, you can write it as $\phi^{i-n} + \Delta x \frac{d\phi}{dx} + \frac{\Delta x^2}{2!} \frac{d^2\phi}{dx^2} - \frac{\Delta x^3}{3!} \frac{d^3\phi}{dx^3} + \dots$

Then you will get $-\Delta x \Delta t \frac{d^2\phi}{dx dt} + \frac{\Delta x^2 \Delta t^2}{2!} \frac{d^2\phi}{dx^2 dt^2} - \frac{\Delta x^3 \Delta t^3}{3!} \frac{d^3\phi}{dx^3 dt^3} + \dots$ So, $\frac{d^3\phi}{dx^3 dt^3} + 3 \frac{\Delta x^2 \Delta t^2}{6} \frac{d^3\phi}{dx^2 dt^2} - 3 \frac{\Delta x \Delta t^2}{6} \frac{d^3\phi}{dx dt^2} + \dots$

So, now let us see that which are terms are cancelled out so one is you can see this ϕ^{i+n} this ϕ^{i-n} and this $\frac{1}{2} \phi^{i+n}$, so these are cancelled out so $\Delta x \frac{d\phi}{dx}$ this is your plus and this is your minus. Then you can see this is $\Delta x \Delta t$ and this is so this term and this term so this you can cancel then $\Delta t \frac{d\phi}{dt}$ and so this is one term, this is one term, this is one term, this is plus 2 and this is your minus. So, this you can cancel.

Then you can see what are the other terms will be cancelled you just check. So, you can see $\frac{\Delta t^2}{2!} \frac{d^2\phi}{dt^2}$ so these, this term, this term and you can see this term. So, now these you can cancel out, because these is minus 2 then you have $\frac{\Delta x^3}{6} \frac{d^3\phi}{dx^3}$ so you can see these and these will cancel out. Then you will get $3 \Delta x \Delta x \Delta t^2$ so this term and this term will get cancel.

Then what are the remaining terms let us see, so you will get the remaining terms is this one, one is this, then you have this term then you have this term you have $\Delta x^2 \frac{d^2\phi}{dx^2}$ so this term. And you can see this term also you will get cancelled out so you can see this is your, this is your 6 actually. So, you can see this term, this term so it is 2 into this term and it is minus 2 into this term so this will get canceled out.

So, now you can see these are the 4 terms remaining from this expansion and that you can write now is equal to gamma by 2 delta x square, first let us write. Then it will be easy, so you can see this delta x square by 2 delta x square by 2 so there are 2 so 2 times delta x square by 2 del 2 phi by del x square. So, this is also 2 times so you can write 2 into 3 so 2 into 3 delta x square delta t by 6 del cube phi by del x square del t.

So, n plus higher order term we are neglecting. So, you can see you will get gamma by 2 so it will be just, so 2 2 will get cancelled out and delta x square so it is del 2 phi by del x square. So, this three 6 6 will get cancelled out so you will get delta t only and del cube phi by del x square and delta t. So, now we have found the left hand side term and term 1 and term 2 in the right hand side, so if you put all together then what you are going to get.

(Refer Slide Time: 61:36)

Parabolic Equations

LHS = Term I of RHS + Term II of RHS

$$\frac{\partial \phi}{\partial t} + \frac{\Delta t}{2} \frac{\partial^2 \phi}{\partial t^2} + \frac{(\Delta t)^2}{6} \frac{\partial^3 \phi}{\partial t^3} + \dots = \frac{\Gamma}{2} \left[\frac{\partial \phi}{\partial x^2} + \Delta t \frac{\partial^3 \phi}{\partial x^2 \partial t} + \dots \right] + \frac{\Gamma}{2} \left[\frac{\partial^2 \phi}{\partial x^4} + \frac{(\Delta x)^2}{12} \frac{\partial^4 \phi}{\partial x^4} + \dots \right]$$

$$\frac{\partial \phi}{\partial t} = \Gamma \frac{\partial^2 \phi}{\partial x^2} \rightarrow \frac{\partial \phi}{\partial t} - \Gamma \frac{\partial^2 \phi}{\partial x^2} = 0$$

$$\frac{\partial^2 \phi}{\partial t^2} = \Gamma \frac{\partial^3 \phi}{\partial x^2 \partial t}$$

$$\frac{\partial \phi}{\partial t} - \Gamma \frac{\partial^2 \phi}{\partial x^2} = \frac{\Delta t}{2} \frac{\partial^2 \phi}{\partial t^2} - \frac{\Delta t}{2} \frac{\partial^3 \phi}{\partial x^2 \partial t} + \frac{\Gamma (\Delta x)^2}{24} \frac{\partial^4 \phi}{\partial x^4} - \frac{(\Delta t)^2}{6} \frac{\partial^3 \phi}{\partial t^3} + \dots$$

Truncation Error

$$TE = \frac{\Gamma}{24} (\Delta x)^2 \frac{\partial^4 \phi}{\partial x^4} - \frac{(\Delta t)^2}{6} \frac{\partial^3 \phi}{\partial t^3} + \dots \text{HOT}$$

$$O[(\Delta x)^2, (\Delta t)^2] - \text{Crank-Nicolson method}$$

So, in the left hand side whatever you got so you will get so we are now writing left hand side is equal to term 2 of right hand side plus term 1 of right hand side so if you write it then what you are going to get? So, left hand side you can see it is there is del phi by del t plus del t by 2 del 2 phi by del t square plus del t square by 6 del cube phi by del t cube.

So, you can see this is the term, we have written, now other 2 terms, term 1 and term 2 whatever we have derived so let us write it equal to so it will be gamma by 2, gamma by 2. So, what is there in the second term gamma by 2 so it is del 2 phi by del x square del 2 phi by del x square plus del t del cube phi by del x square del t.

This is the term 2 and what is term 1? It is also γ by 2 so γ by 2 so term 2 is you see term 2 whatever we have written, so this is the term 2. So, this if you write it here so you are going to get $\frac{\partial^2 \phi}{\partial x^2} + \frac{\Delta x^2}{12} \frac{\partial^4 \phi}{\partial x^4} + \text{high order term}$.

So, now we have written this whatever from the expansion we have got so in the finite difference equation we put and we have written now we will find the truncation error so you can see that our governing equation is $\frac{\partial \phi}{\partial t} = \gamma \frac{\partial^2 \phi}{\partial x^2}$. Let us take the time derivative of this equation, so $\frac{\partial^2 \phi}{\partial t^2}$ you can write as $\gamma \frac{\partial^2 \phi}{\partial x^2} \frac{\partial t}{\partial t}$. So, you can see that whatever term is there here that we can replace with $\frac{\partial^2 \phi}{\partial t^2}$.

So, you replace it there, so what you are going to get? So, $\frac{\partial \phi}{\partial t} = \gamma \frac{\partial^2 \phi}{\partial x^2}$ all you take in the right hand side, so now you can see this is your γ by 2 $\frac{\partial^2 \phi}{\partial x^2}$ and it is γ by 2 $\frac{\partial^2 \phi}{\partial x^2}$. So, it will be just $\gamma \frac{\partial^2 \phi}{\partial x^2}$, so that you can take in the left hand side. So, it will be minus so this is the so you can see that $\gamma \frac{\partial^2 \phi}{\partial x^2} t = \frac{\partial^2 \phi}{\partial t^2}$ so you can write it as $\frac{\partial^2 \phi}{\partial t^2} = \gamma \frac{\partial^2 \phi}{\partial x^2}$.

So, the first term this term we have written here and so is your $\frac{\Delta t}{2}$ because γ by 2 is there, so now you take it in the right hand side so it is v minus $\frac{\Delta t}{2} \frac{\partial^2 \phi}{\partial t^2}$. So, you can see this will get cancel out and this if you take and what is left here you have γ by 2 so it will be plus γ by 24 $\frac{\partial^2 \phi}{\partial x^2}$ to the power 4. This is there and if you take this term in the minus $\frac{\Delta t^2}{6} \frac{\partial^3 \phi}{\partial t^3}$, some other terms.

So, this get cancelled out so you can see that your governing equation is this one, so that means $\frac{\partial \phi}{\partial t} = \gamma \frac{\partial^2 \phi}{\partial x^2}$, so this is equal to 0 but from the finite difference equation we got this additional term so this is your truncation error. So, now what is your truncation error? So, truncation error now so truncation error TE is γ by 24 so this Δx^2 will be there.

Here we miss Δx^2 , so this term is coming here so Δx^2 to the power Δx^2 so it will be $\Delta x^2 \frac{\partial^4 \phi}{\partial x^4} - \frac{\Delta t^2}{6} \frac{\partial^3 \phi}{\partial t^3} + \text{other high order terms}$. So, now you can see the truncation error so the leading term

what is there Δx square and Δt square. So, the leading terms, in the truncation error we have Δx square and Δt square that means the order of accuracy of this Crank-Nicolson scheme is Δx square and Δt square so we have proved it now.

So, the accuracy is order of Δx square and Δt square. So, it is just truncation error we have found which is the difference between the partial differential equation and the finite difference equation. So, first we discretize that equation after that we expand it using Taylor series each dependent variables then we put it in the finite difference equation then we the difference we have found as truncation error and the leading terms you can see that it is Δx square and Δt square.

So, obviously the you can see that Crank-Nicolson method is your, Crank-Nicolson method is Crank-Nicolson method is second order or accurate in both time and space. Thank you.